ON THE HEAT EQUATION INVOLVING A GRUSHIN OPERATOR IN MARCINKIEWCZ SPACES

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ABSTRACT. In this work, we give sufficient conditions to the existence and uniqueness for the heat equation involving the operator

$$\Delta_{\mathcal{G}} = \frac{1}{2} \left(\Delta_x + |x|^2 \Delta_y \right)$$

in Marcinkiewcz spaces. Furthermore, we give sufficient conditions to the existence of positive, symmetric and self-similar solutions.

1. Introduction and main results

Since Fujita's paper [12], the semilinear heat equation

(1.1)
$$u_t(x,t) = \Delta u(x,t) + |u(x,t)|^{\rho-1} u(x,t), \text{ in } (0,\infty) \times \mathbb{R}^N,$$

(1.2)
$$u(x,0) = u_0(x)$$
, in \mathbb{R}^N

 $\rho > 1$, is of great mathematical interest.

It is well-known that if $u_0 \in C_0(\mathbb{R}^N)$ be nonnegative and nonzero, then

- (1) If $1 < \rho < 1 + \frac{2}{N}$, there exists no positive global solution of (1.1)–(1.2).
- (2) If $\rho > 1 + \frac{2}{N}$, there exists $u_0 \in L^{\frac{N}{2}(\rho-1)}(\mathbb{R}^N)$ such that there exists a global positive solution of (1.1)–(1.2).

The critical case $\rho = 1 + \frac{2}{N}$ was resolved by Weissler [28]. Initial conditions in Lebesgue spaces were considered in [4, 28].

In this work, we will replace the classical Laplacian with the operator

$$\Delta_{\mathcal{G}} = \frac{1}{2} \left(\Delta_x + |x|^2 \Delta_y \right),\,$$

where Δ_x, Δ_y denote the classical Laplacian in the variables $x \in \mathbb{R}^N$ and $y \in \mathbb{R}^k$, respectively.

This operator, nowadays called of Grushin-type (see [15, 16]), actually belongs to the general class of the operators studied by Hörmander in [17] and is hypoelliptic. Moreover, it is also a particular case of the class of degenerate elliptic operators studied in [18] by Kogoj and Lanconelli. The first to introduce and study a metric and a consequent underlying sub-Riemannian structure for such operators were Franchi and Lanconelli in their seminal papers in the early 1980s [10, 9, 11]. In recent years,

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many works appeared in the literature dealing with second order linear and semilinear degenerate elliptic PDOs falling into this class. We refer to [1, 2, 7, 19, 22] and reference therein for recent results involving the Grushin-type operators.

We are dealing with the following problem

(1.3)
$$u_t = \Delta_{\mathcal{G}} u + |u|^{\rho-1} u, \text{ in } (0, \infty) \times \mathbb{R}^{N+k},$$
$$u(0) = u_0, \text{ in } \mathbb{R}^{N+k}.$$

In bounded domains, we can find well-posedness, longtime dynamics and the existence of attractors for semilinear equations involving degenerate elliptic operators containing the Grushin ones in [20, 21, 24]. In particular, the authors of [24] applied the Galerkin method to prove the existence of solutions of parabolic and pseudo-parabolic equations with associated with Hörmander-type operators, where they used Sobolev-type estimates, which is available for bounded domains. They also proved the existence and upper continuity of attractors.

In [26], Oliveira and Viana prove existence, uniqueness, continuous dependence and blowup alternative of local mild solutions for (1.3)–(1.4) with initial conditions in Lebesgue spaces. Also, they obtain the existence of global solutions in the special case of $u_0 \in L^{\frac{N}{2}(\rho-1)}(\mathbb{R}^{N+k})$ with the sufficiently small norm.

In this manuscript, by working in Marcinkiewcz spaces, we allow larger (in Lebesgue norm) initial data to be taken into account in order to obtain global solutions. More precisely, we give sufficient conditions to the existence and uniqueness of mild solutions for (1.3)–(1.4), with initial conditions in the critical Marcinkiewcz space $L^{p,\infty}(\mathbb{R}^{N+k})$, with

(1.4)
$$p = \frac{N+2k}{2}(\rho - 1).$$

Then, we prove the existence of positive, symmetric and self-similar solutions. Indeed, for example, if $u_0(x,y) = \varepsilon |x|^{-\frac{2}{\rho-1}}|y|^{-\frac{1}{\rho-1}}$ and $\varepsilon > 0$ is sufficiently small, Theorems 1.1 (a) and 1.2 gives the existence of a positive, self-similar and symmetric solution of (1.3)–(1.4) in $L^{\infty}((0,\infty);L^{(p,\infty)})$. Notice that such an initial condition has infinity $L^p(\mathbb{R}^{N+k})$ -norm so we cannot apply those results in [26]. In other words, Marcinkiewcz spaces allow singular homogeneous initial conditions that generate positive, self-similar solutions. Moreover, Theorem 1.1 also gives sufficient conditions for uniqueness, regularity and time-decay of the mild solutions.

The number N+2k is the so-called homogeneous dimension attached to our operator (see [19]).

Here, we will rely on the *explicit* expression of the heat kernel of the Grushin operator (see e.g. [26, 13]): (1.5)

$$K(x, x_0, y; t) = \frac{1}{(2\pi)^{\frac{N+2k}{2}}} \int_{\mathbb{R}^k} \left(\frac{|\xi|}{\sinh(|\xi|t)} \right)^{\frac{N}{2}} e^{i\xi \cdot y - \frac{|\xi|}{2} \left((|x|^2 + |x_0|^2) \coth(|\xi|t) - 2x \cdot x_0 \operatorname{csch}(|\xi|t) \right)} d\xi,$$

for
$$(x, x_0, y) \in \mathbb{R}^{2N+k}, t > 0$$
.

To prove this, the authors followed the Geometric Method in [5] to find the heat kernel of the heat equation with quadratic potential and then applied the partial Fourier transform on the variable y. Then, the inverse partial Fourier transform gave (1.5), an expression that also appears in Theorem 3.4 of the paper by Garofalo and Tralli [13], with slight differences due to constant choices in the definition of the operator.

Let us recall some properties of the heat kernel and the heat semigroup in L^p proved in [26]. The kernel (1.5) is C^{∞} and

- (1) $K_t = \Delta_{\mathcal{G}} K$;
- (2) $\int_{\mathbb{R}^{N+k}} K(x,0,y;t) d(x,y) = 1$;
- (3) $\lim_{t\to 0^+} \int_{\mathbb{R}^{N+k}} K(x, w, y z; t) \varphi(w, z) d(w, z) = \varphi(x, y)$.

Moreover,

(1.6)
$$K(x, x_0, y; t) = t^{-\frac{N+2k}{2}} K(t^{-\frac{1}{2}}x, t^{-\frac{1}{2}}x_0, t^{-1}y; 1).$$

and the heat kernel (1.5) is positive. In particular, $||K(\cdot,0,\cdot;t)||_{L^1(\mathbb{R}^{(N+k)})} = 1$. The unique solution of the Cauchy problem

(1.7)
$$\partial_t u - \Delta_{\mathcal{G}} u = 0, \quad x \in \mathbb{R}^{N+k}, \ t > 0$$

(1.8)
$$u(x, y, 0) = u_0(x, y), \qquad x \in \mathbb{R}^{N+k},$$

defines a strongly continuous semigroup in L^p . It is defined by

(1.9)
$$S_{\mathcal{G}}(t)\varphi(x,y) = \int_{\mathbb{R}^{N+k}} K(x,w,y-z;t)\varphi(w,z) d(w,z).$$

Indeed, we have the following result proved in [26].

Theorem A. For all $1 \leq p \leq \infty$, $S_{\mathcal{G}}(t) : L^p(\mathbb{R}^{N+k}) \to L^p(\mathbb{R}^{N+k})$ is a semigroup. If $1 \leq p \leq r \leq \infty$, then

(1.10)
$$||S_{\mathcal{G}}(t)\varphi||_{L^{r}(\mathbb{R}^{N+k})} \leq C||\varphi||_{L^{p}(\mathbb{R}^{N+k})} t^{-\frac{N+2k}{2}\left(\frac{1}{p}-\frac{1}{r}\right)};$$

Furthermore, for $1 \le p \le r \le \infty$, for all $t_0 > 0$, it is strongly continuous, that is, for $\varphi \in L^r(\mathbb{R}^{N+k})$, it holds

(1.11)
$$||S_{\mathcal{G}}(t)\varphi - S_{\mathcal{G}}(t_0)\varphi||_{L^p} \to 0,$$

as $t \to t_0$. When $p = r < \infty$, then we may take $t_0 = 0$.

Now, we state our results. In the following, $X = L^{\infty}((0, \infty); L^{(p, \infty)}(\mathbb{R}^{N+k}))$ and a mild solution is a solution in X that satisfies the integral equation associated with (1.3)–(1.4) by means of the Duhamel principle. We refer to Section 3 for the precise definition of the mild solution definition.

Theorem 1.1. Let 1 be given by (1.4).

(a) **Well-posedness.** There exists $\delta_p > 0$ sufficiently small such that , if $\|u_0\|_{(p,\infty)} < \delta_p$, then the problem (1.3)-(1.4) has a global mild solution $u \in X$ that is unique in $B_X(2\varepsilon)$ ($\varepsilon > 0$ will be precise later). Furthermore, if $u, v \in are$ mild solutions of (1.3)-(1.4) with initial conditions u_0, v_0 with $L^{(p,\infty)}$ norm less than δ_p , respectively, then

$$||u(t) - v(t)||_{(p,\infty)} \le \frac{1}{1 - 2^{\rho} K_n \varepsilon^{\rho - 1}} ||u_0 - v_0||_{(p,\infty)}$$

- (b) **Regularity.** If $u_0 \in L^{(p,\infty)}(\mathbb{R}^{N+k}) \cap L^{(q,\infty)}(\mathbb{R}^{N+k})$, $||u_0||_{(q,\infty)} < \delta_q$, with $\delta_q > 0$ sufficiently small, for $q > \frac{N+2k}{N+2(k-1)}$, then $u \in L^{\infty}((0,\infty); L^{(p,\infty)}(\mathbb{R}^{N+k}) \cap \mathbb{R}^{N+k})$ $L^{(q,\infty)}(\mathbb{R}^{N+k})$. Moreover, for p < r < q, we have $u, u_t \in L^{\infty}(0,\infty; L^r(\mathbb{R}^{N+k}))$ and $u(t,\cdot) \in C^{\infty}(\mathbb{R}^{N+k})$, for each t > 0.
- (c) **Decay.** Moreover, for $q > \rho p$, we have

$$||u(t)||_{(r,\infty)} \le Ct^{-\sigma}, \ t > 0.$$

(d) Uniqueness. The solution in unique in the set $L^{\infty}((0,\infty);L^{(p,\infty)}(\mathbb{R}^{N+k})\cap$ $L^{(q,\infty)}(\mathbb{R}^{N+k})$), for 1 .

One key ingredient for the proof of Theorem 1.1 is a Yamazaki-type inequality ([29]): if 1 , there is a constant <math>C > 0 such that

(1.14)
$$\int_0^\infty t^{\frac{n}{2}\left(\frac{1}{p}-\frac{1}{q}\right)-1} \|G(t)\phi\|_{(q,1)} ds \le C \|\phi\|_{(p,1)},$$

for each $\phi \in L^{(p,1)}(\mathbb{R}^n)$, in which G(t) denotes the heat semigroup. Also, it was generalized by Ferreira and Villamizar-Roa, given by (see [8]) for the semigroup generated by the fractional Laplacian. We will need to use a similar inequality for the semigroup generated by our operator (see Lemma 2.5 below). The regularity relies on interpolation arguments and on the smoothness of the kernel.

In the next theorem, we will use the following terminology. By a self-similar solution, we mean a solution u that satisfies the scaling map

(1.15)
$$u_{\lambda}(x,y,t) := \lambda^{\frac{2}{\rho-1}} u(\lambda x, \lambda^2 y, \lambda t).$$

Let \mathcal{A} be a subset of the orthogonal matrices in $\mathcal{O}(N) \times \mathcal{O}(k)$. We say that a function φ is invariant under the action of \mathcal{A} if $\varphi(Tz) = \varphi(z)$, for all $z \in \mathbb{R}^{N+k}$ and $T \in \mathcal{A}$.

Theorem 1.2. Let the assumptions of Theorem 1.1 hold.

- (a) If u_0 is a nonnegative function, then u is positive.
- (b) If u₀ is a symmetric function under the action of A, then u also is.
 (c) If u₀ is a homogeneous function ²/_{\rho-1}, then the solution u given by Theorem 1.1 is self-similar.

The rest of the manuscript is organized as follows. In Section 2, we define and gather some properties of the Lorentz and Marcinkiewcz spaces. We also study properties of the semigroup associated with our operator in Lorentz spaces, and the Yamazaki-type estimate. Other key estimates are proved in this section. Section 3 is devoted to the proof of Theorem 1.1, and Section 4 to the proof of the symmetries. We close the paper with some remarks on the solutions to our problem.

2. Key results

Most of the notation we use in this paper is standard. Lorentz and Marcinkiewicz spaces $L^{(p,q)}$ and $L^{(p,\infty)}$ are as defined in [3, 8]. The norm in these spaces will be denoted by $\|\cdot\|_{(p,q)}$. Next, we collect the main properties of these spaces.

Consider a measurable function $f: \mathbb{R}^n \to \mathbb{R}$. The f^* rearrangement function is defined by

$$f^*(t) = \inf\{s > 0 : m(\{x \in \mathbb{R}^n : |f(x)| > s\}) \le t\}, \ t > 0,$$

where m is the measure \mathbb{R}^n -Lebesgue. The averaging function f^{**} is defined by

$$f^{**}(t) = \frac{1}{t} \int_0^t f^*(s) ds, \ t > 0.$$

Let $0 . The Lorentz space, <math>L^{(p,q)}(\mathbb{R}^n)$, is the set of all measurable functions $f: \mathbb{R}^n \longrightarrow \mathbb{R}$, such that $||f||_{(p,q)}^* < \infty$, where

(2.1)
$$||f||_{(p,q)}^* = \begin{cases} \left[\frac{q}{p} \int_0^\infty \left(t^{\frac{1}{p}} f^*(t) \right)^q \frac{dt}{t} \right]^{\frac{1}{q}}, \ t > 0 & \text{if } 0 0} t^{\frac{1}{p}} f^*(t), & \text{if } 0$$

Since $||f||_{(p,q)}^*$ does not satisfy the triangular inequality, $L^{(p,q)}$ is metrizable with the norm $||f||_{(p,q)}$ given by

$$(2.2) ||f||_{(p,q)} = \begin{cases} \left(\frac{q}{p} \int_0^\infty \left[t^{\frac{1}{p}} f^{**}(t)\right]^q dt/t\right)^{\frac{1}{q}}, & \text{if } 1 0} t^{\frac{1}{p}} f^{**}(t), & \text{if } 1$$

When $q = \infty$, the space $L^{(p,\infty)}$ is called the Marcinkiewics space or weak- L^p .

Proposition 2.1. The spaces $L^{(p,q)}$ with the norms $||f||_{(p,q)}$ are Banach spaces and

(2.3)
$$||f||_{(p,q)}^* \le ||f||_{(p,q)} \le \frac{p}{p-1} ||f||_{(p,q)}^*,$$

where $1 and <math>1 \le q \le \infty$. If, in addition, (\mathbb{R}^n, μ) be a σ -finite space, then

$$\begin{array}{rcl} (L^{(p,q)})^* & = & L^{(p,\infty)}, & & 1$$

Moreover, given $T \in (L^{(p,q)})^*$, with $1 and <math>0 < q \le \infty$, exists g a measurable function such that

(2.4)
$$T(f) = \int_{\mathbb{R}^n} fg d\mu.$$

The following remark is a consequence of the equivalence between the norm and the seminorm of the Lorentz space.

Remark 2.2. Let $1 and <math>0 < q < \infty$. If $h \in L^{(p,\infty)}(\mathbb{R}^n)$, then $||h|^q||_{(p,\infty)} \leq \frac{p}{p-1} ||h||^q_{(pq,\infty)}$.

Plus, we need a Hölder-type estimate.

Proposition 2.3 (Generalized Hölder inequality). Let $1 < p_1, p_2 < \infty$. Let $f \in L^{(p_1,d_1)}$, $g \in L^{(p_2,d_2)}$, and $\frac{1}{p_1} + \frac{1}{p_2} < 1$. Then the product h = fg belongs to $L^{(r,d_3)}$, where $\frac{1}{r} = \frac{1}{p_1} + \frac{1}{p_2}$, and $d_3 \ge 1$ is such that $\frac{1}{d_1} + \frac{1}{d_2} \ge \frac{1}{d_3}$. Moreover,

$$||h||_{(r,d_3)} \le C(r)||f||_{(p_1,d_1)}||g||_{(p_2,d_2)}.$$

If $r = d_3 = 1$, then $h \in L^1$ and

$$||h||_1 \le ||f||_{(p,q_1)} ||g||_{(p',q_2)},$$

where p' is the conjugate of p.

We can use the interpolation theory in [3] and Theorem A to obtain the behavior of the semigroup generated by our operator in Lorentz spaces.

Proposition 2.4. For all $1 \leq p \leq \infty$, the semigroup $S_{\mathcal{G}}(t) : L^{(p,s)}(\mathbb{R}^{N+k}) \to L^{(r,s)}(\mathbb{R}^{N+k})$ is a strongly continuous semigroup for t > 0, provided that $1 \leq p \leq r \leq \infty$ and $1 < s \leq \infty$. Moreover,

Moreover, given $1 and <math>\varphi \in L^{(p',1)}$, we have

(2.6)
$$\langle S_{\mathcal{G}}(t)v, \varphi \rangle = \langle v, S_{\mathcal{G}}(t)\varphi \rangle.$$

Proof. Recall that Lorentz spaces are real interpolation of Lebesgue spaces, that is, $L^{p,q} = (L^{p_0}, L^{p_1})_{\eta,q}$ with $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$. Also, Th. A gives that $S_{\mathcal{G}}: L^{p_i}(\mathbb{R}^{N+k}) \to L^{r_i}(\mathbb{R}^{N+k})$, $1 \leq p_i \leq r_i \leq \infty$, i = 0, 1. From the Marcinkiewcz interpolation theorem (see [3]) and Theorem A, we obtain that

$$\begin{split} \|S_{\mathcal{G}}(t)\|_{L^{(p,q)} \to L^{(r,q)}} &\leq c \|S_{\mathcal{G}}(t)\|_{L^{p_0} \to L^{r_0}}^{1-\eta} \|S_{\mathcal{G}}(t)\|_{L^{p_1} \to L^{r_1}}^{\eta} \\ &\leq C \left(t^{-\frac{N+2k}{2}\left(\frac{1}{p_0} - \frac{1}{r_0}\right)}\right)^{1-\eta} \left(t^{-\frac{N+2k}{2}\left(\frac{1}{p_1} - \frac{1}{r_1}\right)}\right)^{\eta} \\ &= Ct^{-\frac{N+2k}{2}\left(\frac{1}{p} - \frac{1}{r}\right)}, \end{split}$$

because $\frac{1}{p} = \frac{1-\eta}{p_0} + \frac{\eta}{p_1}$ and $\frac{1}{r} = \frac{1-\eta}{r_0} + \frac{\eta}{r_1}$.

Identity (2.6) comes from a combination of the duality relation $(L^{(p,\infty)})^* = L^{(p',1)}$ and Fubini's theorem.

Now, let us prove a Yamazaki-type lemma (see [29]).

Lemma 2.5. Let 1 . Then,

(2.7)
$$\int_{0}^{\infty} t^{\frac{N+2K}{2}\left(\frac{1}{p}-\frac{1}{r}\right)-1} \|S_{\mathcal{G}}(t)\phi\|_{(r,1)} ds \le c \|\phi\|_{(p,1)}$$

Proof. Set $\xi(t) = t^{\frac{N+2k}{2}(\frac{1}{p}-\frac{1}{r})-1} \|S_{\mathcal{G}}(t)\phi\|_{(r,1)}$ and $1 < p_1 < p < p_2 < r$ such that

$$\frac{1}{p} = \frac{\lambda}{p_1} + \frac{1-\lambda}{p_2}, \lambda \in (0,1).$$

From Proposition 2.4, we have

$$||S_{\mathcal{G}}(t)\phi||_{(r,1)} \le ct^{-\frac{N+2k}{2}\left(\frac{1}{p_j}-\frac{1}{r}\right)}||\phi||_{(p_j,1)}, \ \phi \in {}^{(p_j,1)}(\mathbb{R}^{N+k}), \ j=1,2.$$

Now, define $\frac{1}{l_j} = \frac{N+2K}{2} \left(\frac{1}{p_j} - \frac{1}{p} \right) + 1$. The above estimate gives us

$$\xi(t) \le c \ t^{l_j^{-1}} \|\phi\|_{(p_j,1)}.$$

Since $||t^{l_j^{-1}}||_{(p_j,\infty)} = 1$, we conclude that $||\xi(t)||_{(l_j,\infty)} \leq c||\phi||_{(p_j,1)}$. Note that $1 = \frac{\lambda}{l_1} + \frac{1-\lambda}{l_2}$, $L^{(p,1)}(\mathbb{R}^{N+k}) = (L^{(p_1,1)}(\mathbb{R}^{N+k}), L^{(p_2,1)}(\mathbb{R}^{N+k}))_{\lambda,1}$ and $L^{(1,1)}(0,\infty) = (L^{(l_1,0)}(0,\infty), L^{(l_2,0)}(0,\infty))_{\lambda,1}$. Then, real interpolation again gives us $\xi(t)$: $L^{(p,1)}(\mathbb{R}^{N+k}) \to L^{(1,1)}(0,\infty)$ and

$$\|\xi(t)\|_{(1,\infty)} \le c \|\phi\|_{(p,1)},$$

that is (2.7).

For the function f be the function $f(w) = |w|^{\rho-1}w, w \in \mathbb{R}, 1 < \rho < r < \infty$, and we define

$$\zeta(h)(x) = \int_0^\infty S_{\mathcal{G}}(s)(f(h))(s) \ ds.$$

We have the following lemma.

Lemma 2.6. Let
$$h \in L^{\infty}\left((0,\infty); L^{(p,\infty)}\right)$$
, $p = \frac{N+2k}{2}(\rho-1)$, and $\rho > \frac{N+2k}{N+2(k-1)}$. Then, (2.8) $\|\zeta(h)\|_{(p,\infty)} \le K \sup_{t>0} \|h(t)\|_{(p,\infty)}^{\rho}$.

If, in addition, $h \in L^{\infty}\left((0,\infty); L^{(p,\infty)} \cap L^{(q,\infty)}\right)$, for $q > \frac{N+2k}{N+2(k-1)}$, then

(2.9)
$$\|\mathcal{C}(h)\|_{(q,\infty)} \le C \sup_{t>0} \|h(t)\|_{(p,\infty)}^{\rho-1} \sup_{t>0} \|h(t)\|_{(q,\infty)}.$$

Proof. Let $\phi \in L^{(p,\infty)}$, note that $\frac{p}{p-\rho} > p' = \frac{p}{p-1}$ and $\frac{N+2k}{2} \left(\frac{1}{p'} - \frac{p-\rho}{p}\right) - 1 = 0$. From (2.6), the Hölder inequality 2.3, and the Yamazaki-type inequality (2.7), we have

$$\begin{split} \|\zeta(h)\|_{(p,\infty)} &= \sup_{\|\phi\|_{(p',1)} = 1} \left| \int_{\mathbb{R}^{N+K}} \zeta(h) \phi \right| \\ &= \sup_{\|\phi\|_{(p',1)} = 1} \left| \int_{0}^{\infty} \int_{\mathbb{R}^{N+k}} (S_{\mathcal{G}}(s)f(h))(s) \phi(z) \ dz \ ds \right| \\ &= \sup_{\|\phi\|_{(p',1)} = 1} \left| \int_{0}^{\infty} \int_{\mathbb{R}^{N+k}} f(h)(S_{\mathcal{G}}(s)\phi) \ dz \ ds \right| \\ &\leq \sup_{\|\phi\|_{(p',1)} = 1} \int_{0}^{\infty} \|f(h)\|_{\left(\frac{p}{\rho},\infty\right)} \|S_{\mathcal{G}}(s)\phi\|_{\left(\frac{p}{\rho-\rho},1\right)} \ ds \\ &\leq C \sup_{\|\phi\|_{(p',1)} = 1} \int_{0}^{\infty} t^{\frac{N+2k}{2}\left(\frac{1}{p'} - \frac{p-\rho}{p}\right) - 1} \|S_{\mathcal{G}}(s)\phi\|_{\left(\frac{p}{p-\rho},1\right)} ds \left(\sup_{t > 0} \|h(t)\|_{(p,\infty)}^{\rho}\right) \\ &\leq C \sup_{t > 0} \|h(t)\|_{(p,\infty)}^{\rho}. \end{split}$$

This proves (2.6).

To prove (2.9), we repeat the idea above rather using that

$$\|\mathcal{C}(h)\|_{(q,\infty)} \le \sup_{\|\phi\|_{(r',1)}=1} \int_0^\infty \|f(h)\|_{(r,\infty)} \|S_{\mathcal{G}}(t)\phi\|_{(r',1)} dt$$

and

$$||f(h)||_{(r,\infty)} \le ||h||_{(p,\infty)}^{\rho-1} ||h||_{(q,\infty)}$$

for $\frac{1}{r} = \frac{\rho-1}{p} + \frac{1}{q}$. Now, we apply Yamazaki' inequality with (p,q) = (q',r') to obtain

$$\|\mathcal{C}(h)\|_{(q,\infty)} \le C \sup_{t>0} \|h(t)\|_{(p,\infty)}^{\rho-1} \sup_{t>0} \|h(t)\|_{(q,\infty)},$$

since
$$q' < r'$$
, and $\frac{N+2K}{2} \left(\frac{1}{q'} - \frac{1}{r'} \right) = \frac{N+2K}{2} \left(\frac{1}{r} - \frac{1}{q} \right) = \frac{N+2K}{2} \frac{(\rho-1)}{p} = 1$.

Let us recall the following contraction-principle based in [8, Lemma 3.9]. It is especially useful to control the size of the constants that will appear in some proofs.

Lemma 2.7. Let $1 < \rho < \infty$ and X be a Banach space with norm $\|\cdot\|$, and $B: X \to X$ be a map witch satisfies

$$||B(x)|| \le K||x||^{\rho}$$

and

$$||B(x) - B(z)|| \le K||x - z|| (||x||^{\rho - 1} + ||z||^{\rho - 1}).$$

Let R > 0 be the unique positive root of the equation $2^{\rho}Ka^{\rho-1} - 1 = 0$. Given $0 < \varepsilon < R$ and $y \in X, y \neq 0$, such that $||y|| < \varepsilon$, there exists a solution $x \in X$ for the equation x = y + B(x) such that $||x|| \leq 2\varepsilon$. The solution x is unique in the ball $B_{2\varepsilon} := \bar{B}(0, 2\varepsilon)$. Moreover, the solution depends continuously on y in the following sense: If $||\tilde{y}|| \leq \varepsilon, \tilde{x} = \tilde{y} + B(\tilde{x})$, and $||\tilde{x}|| \leq 2\varepsilon$, then

$$||x - \tilde{x}|| \le \frac{1}{1 - 2\rho K \varepsilon^{\rho - 1}} ||y - \tilde{y}||.$$

3. Proof of Theorem 1.1

This section is devoted to the proof of the global existence of solutions for the problem (1.3)-(1.4) in the space X that will be defined below. As it is claimed in the introduction, we will take the initial data in $L^{(p,\infty)}(\mathbb{R}^{N+2k})$, with $p = \frac{N+2k}{2}(\rho-1)$.

Definition 3.1. We define the space X, formed by the functions $u:(0,\infty)\to L^{(p,\infty)}$, such that

$$u \in L^{\infty}((0,\infty); L^{(p,\infty)}),$$

with the norm $||u||_X = \sup_{t>0} ||u(t)||_{(p,\infty)}$.

Considering this space and the Duhamel principle applied to (1.3)–(1.4), we define mild solution below.

Definition 3.2. For $u_0 \in L^{(p,\infty)}$, a global mild solution of the initial value problem (1.3) - (1.4) is a solution $u \in X$ of the integral equation

(3.1)
$$u(t) = S_{\mathcal{G}}(t)u_0 + \int_0^t S_{\mathcal{G}}(t-s)|u|^{\rho-1}u(s) \ ds := S_{\mathcal{G}}(t)u_0 + B(u)(t),$$

such that $u(t) \to u_0$, as $t \to 0^+$, in the sense of the distributions.

From Proposition 2.3 and Remark 2.2, we have the following simple and useful Lipschitz property.

Lemma 3.3. Let $u, v \in X$ and f given by $f(w) = |w|^{\rho-1}w, w \in \mathbb{R}$, for $1 < \rho < r < \infty$. Then,

(3.2)
$$|||f(u) - f(v)|||_{(\frac{r}{\rho},\infty)} \le \rho ||u - v||_{(r,\infty)} \left[||u||_{(r,\infty)}^{\rho-1} + ||v||_{(r,\infty)}^{\rho-1} \right].$$

Proof of Theorem 1.1 (a). We recall (3.1) and define $\Psi: X \to X$ by

$$\Psi(u)(t) = S_{\mathcal{G}}(t)u_0 + B(u)(t).$$

Let $B_X(R)$ be the closed ball for radius R > 0 as in Lemma 2.7, and centered at the origin: $B_X(R) := \{w \in X : ||w||_X \le R\}$. We will apply Lemma 2.7.

Initially, we see that it follows from Proposition 2.4 that there is C' > 0 such that

$$||S_{\mathcal{G}}(t)u_0||_{(p,\infty)} \le C' ||u_0||_{(p,\infty)},$$

for all t > 0. For a sufficiently small $\delta_n > 0$, we conclude that

$$(3.3) ||S_{\mathcal{G}}(t)u_0||_{(p,\infty)} \le \varepsilon,$$

where $0 < \varepsilon < R$.

Now, as in [8], we define $h(s,\cdot) = u(t-s,\cdot)$, if $0 \le s \le t$, and $h(s,\cdot) = 0$ otherwise. Then, $B(u) = \zeta(h)$. Lemma 2.6 then yields

$$(3.4) ||B(u)(t)||_{(p,\infty)} \le K_p \left(||u||_{(p,\infty)}\right)^{\rho},$$

for some $K_p > 0$. At this point, notice that R > 0 chosen as in Lemma 2.7 depends on this $K_p > 0$.

Analogously, for $u, v \in B_X(R)$, we can use (3.2) and Lemma 2.6 to estimate

We are ready to apply Lemma 2.7, which gives us a unique $u \in B_X(2\varepsilon)$ that is a solution of the integral equation $u = \Psi(u)$.

It remains to show that the solution $u(t) \to u_0$, as $t \to 0^+$, in the sense of the distributions. Let $\varphi \in C_0^{\infty}(\mathbb{R}^{N+k})$. Since $C_0^{\infty} \subset L^{(p',1)}$ (not densely injected), then

$$|\langle S_{\mathcal{G}}(t)u_0 - u_0, \varphi \rangle| = |\langle u_0, S_{\mathcal{G}}(t)\varphi - \varphi \rangle| \le ||u_0||_{(p,\infty)} ||S_{\mathcal{G}}(t)\varphi - \varphi||_{(p',1)} \to 0.$$

Similarly, by Fubini's theorem and Hölder inequality

$$|\langle B(u), \varphi \rangle| \le \int_0^t ||u|^{\rho - 1} u||_{(\frac{p}{\rho}, \infty)} ds ||S_{\mathcal{G}}(t - s)\varphi||_{(\frac{p}{p - \rho}, 1)} \le Ct ||\varphi||_{\frac{p}{p - \rho}, 1} \sup_{t > 0} ||u||_r^{\rho} \to 0,$$

as $t \to 0^+$. Therefore, $u(t) \to u_0$, as $t \to 0^+$, in the sense of the distributions.

Furthermore, the last inequality Lemma 2.7 guarantees that, if $u, v \in$ are mild solutions of (1.3)–(1.4) with initial conditions u_0, v_0 , respectively, then

$$||u(t) - v(t)||_{(p,\infty)} \le \frac{1}{1 - 2^{\rho} K_{\nu} \varepsilon^{\rho - 1}} ||u_0 - v_0||_{(p,\infty)}.$$

Proof of Theorem 1.1 (b). Let $u_0 \in L^{(p,\infty)} \cap L^{(q,\infty)}$ implies $u \in L^{\infty}(0,\infty;L^{(q,\infty)})$. Define Picard's sequence

$$u_1 = S_{\mathcal{G}}(t)u_0$$
, and $u_{k+1} = u_1 + B(u_k)$.

We know that $u_k \to u$ in $L^{\infty}(0, \infty; L^{(p,\infty)})$,

$$||u_1||_{(q,\infty)} \le ||u_0||_{(q,\infty)},$$

and

$$||u_{k+1}||_{(q,\infty)} \le C||u_0||_{(q,\infty)} + ||B(u_k)||_{(q,\infty)}.$$

By the part (a), we have that, for all $k \in \mathbb{N}$,

$$||u_{k+1}||_{(p,\infty)} \le 2\varepsilon.$$

From (2.9),

$$||B(u_{k+1})||_{(q,\infty)} \le K_p \sup_{t>0} ||u_k(t)||_{(p,\infty)}^{\rho-1} \sup_{t>0} ||u_k(t)||_{(q,\infty)} \le K_p (2\varepsilon)^{\rho-1} ||u_k(t)||_{(q,\infty)}.$$

for $||u_0||_{(p,\infty)} < \delta_p$. For a possible smaller δ_p , we have $a := K_p(2\varepsilon)^{\rho-1} < 1$. Then

$$||u_{k+1}||_{(q,\infty)} \leq C\delta_q + K_p^{\rho-1} ||u_k||_{(q,\infty)}$$

$$\leq C\delta_q (1 + a + a^2 + \dots + a^k)$$

$$C^{\varsigma}$$

$$= \frac{C\delta_q}{1-a}.$$

Let us see that $\{u_k\}$ is a Cauchy sequence in $L^{(q,\infty)}$. Define $v_k = u_{k+1} - u_k$, Then, (2.9) again will yield

$$||v_{k}||_{(q,\infty)} = ||B(u_{k}) - B(u_{k-1})||_{(q,\infty)}$$

$$= \left\| \int_{0}^{\infty} S_{\mathcal{G}}(s) [f(h_{k}) - f(h_{k-1})] ds \right\|_{(q,\infty)}$$

$$\leq K_{p} \sup_{t>0} ||f(u_{k}) - f(u_{k-1})||_{(r,\infty)}$$

$$\leq K_{p} \sup_{t>0} \left(||u_{k}||_{(p,\infty)}^{\rho-1} + ||u_{k-1}||_{(p,\infty)}^{\rho-1} \right) \sup_{t>0} ||v_{k-1}||_{(q,\infty)}.$$

Now, if $\delta_q > 0$ is sufficiently small, the above estimate and (3.7) ensure that $||v_k||_{(q,\infty)} \to 0$, as $k \to \infty$. Therefore, $\{u_k\}$ converges in $L^{(q,\infty)}$ to a function that equals u by the uniqueness of the limit in the sense of distributions. This proves that $u \in L^{\infty}((0,\infty); L^{(p,\infty)}(\mathbb{R}^{N+k}) \cap L^{(q,\infty)}(\mathbb{R}^{N+k}))$.

In this case, by interpolation,

$$(3.8) t \mapsto u(t) \in L^{\infty}(0, \infty, L^{r}(\mathbb{R}^{N+K}))$$

with p < r < q, because setting η such that $\frac{1}{r} = \frac{\eta}{p} + \frac{1-\eta}{q}$, $q_0, q_1 = \infty$ and s = r, we have

$$(L^{(p,\infty)}, (L^{(q,\infty)})_{(p,r)} = L^{(r,r)} = L^r.$$

This with the smoothness of the heat kernel yield that u(t) is $C^{\infty}(\mathbb{R}^{N+K})$. Moreover, Metafune, Negro and Spina [25] proved that the Grushin operator generates an analytic semigroup $(e^{-\Delta gt})_{t\geq 0}$ in Lebesgue spaces. By uniqueness $e^{-\Delta gt} = S_{\mathcal{G}}(t)$, for all $t\geq 0$,

whence, $(S_{\mathcal{G}}(t)_{t\geq 0})$ is extended to an analytic semigroup. Then, since $u_0 \in L^r(\mathbb{R}^{N+K})$ by interpolation,

(3.9)
$$\left\| \frac{S_{\mathcal{G}}(h) - I}{h} S_{\mathcal{G}}(t) u_0 \right\|_r \to 0,$$

as $h \to 0^+$, for all t > 0. Now, we use the integral representation of the solution, the decay of the semigroup in L^r , and (3.8) to get

$$\left\| \frac{u(t+h) - u(t)}{h} \right\|_{r} \le \left\| \frac{S_{\mathcal{G}}(h) - I}{h} S_{\mathcal{G}}(t) u_{0} \right\|_{r} + 2C \int_{0}^{t} s^{-\frac{p}{r}} \left\| \frac{u(t+h-s) - u(t-s)}{h} \right\|_{r} ds \cdot \sup_{t>0} \|u(t)\|_{r}^{\rho-1}$$

Put
$$\xi_h(t) = \left\| \frac{u(t+h) - u(t)}{h} \right\|_r$$
. Then,

$$\xi_h(t) \le \left\| \frac{S_{\mathcal{G}}(h) - I}{h} S_{\mathcal{G}}(t) u_0 \right\|_{t > 0} + 2C \sup_{t > 0} \|u(t)\|_r^{\rho - 1} \int_0^t s^{-\frac{p}{r}} \xi_h(t - s) \ ds.$$

Because p < r we can apply the Singular Gronwall Lemma (see [4]) to obtain

(3.10)
$$\xi_h(t) \le C \left\| \frac{S_{\mathcal{G}}(h) - I}{h} S_{\mathcal{G}}(t) u_0 \right\|_{\mathcal{F}} \to 0,$$

as
$$h \to 0^+$$
, by (3.9).

Proof of Theorem 1.1 (c). For the decay, it needs $1 , <math>\beta = \frac{N+2k}{2} \left(\frac{1}{p} - \frac{1}{r}\right)$, and the space Y, formed by the functions $u \in X$, such that

$$t \mapsto t^{\sigma} u(t) \in L^{\infty}((0,\infty); L^{(q,\infty)}),$$

endowed with the norm

(3.11)
$$||u||_Y = \sup_{t>0} ||u(t)||_{(p,\infty)} + \sup_{t>0} t^{\beta} ||u(t)||_{(q,\infty)}.$$

In view of the estimates obtained in the proof of part (a), to apply Lemma 2.7 with the Banach space Y, we only need to estimate $t^{\sigma}u(t)$ in $L^{(q,\infty)}$). From Proposition 2.4, we have

$$(3.12) ||S_{\mathcal{G}}(t)u_0||_{(q,\infty)} \le \varepsilon,$$

and (2.9) yields

$$||B(u)||_{(q,\infty)} \le K_p ||u(t)||_{(p,\infty)} ||u(t)||_{(q,\infty)},$$

or

$$\sup_{t>0} t^{\beta} \|B(u)\|_{(q,\infty)} \le K_p \sup_{t>0} \|u(t)\|_{(p,\infty)} \sup_{t>0} t^{\beta} \|u(t)\|_{(q,\infty)}^{\rho-1} \le \|u\|_Y^{\rho}.$$

Then, one can apply Lemma 2.7 to find a mild solution of (1.3)–(1.4) in the ball $B_Y(2\varepsilon)$. In particular,

$$||u(t)||_{(a,\infty)} \le Ct^{-\beta}, \ t > 0.$$

Proof of Theorem 1.1 (d). For the uniqueness, take two mild solutions $u, v \in L^{\infty}((0,\infty); L^{(p,\infty)}(\mathbb{R}^{N+k}) \cap L^{(q,\infty)}(\mathbb{R}^{N+k}))$ for the same Cauchy problem (1.3)–(1.4), with 1 . Then,

$$||u(t) - v(t)||_{(p,\infty)} \leq \int_0^t ||S_{\mathcal{G}}(t-s)[f(u(s)) - f(v(s))]||_{(p,\infty)} ds$$

$$\leq 2C \left(\sup_{t>0} ||u(t)||_{(q,\infty)} \right)^{\rho-1}$$

$$\times \int_0^t (t-s)^{-\frac{N+2k}{2}(\frac{1}{r} - \frac{1}{p})} ||u(s) - v(s)||_{(p,\infty)} ds,$$

with $\frac{1}{r} = \frac{\rho-1}{q} + \frac{1}{p}$. Then, $\frac{N+2k}{2}(\frac{1}{r} - \frac{1}{p}) = \frac{p}{q} < 1$. Define $\varphi(t) = \sup_{s \in (0,t)} \|u(s) - v(s)\|_{(p,\infty)}$. By the Singular Gronwall Lemma, $\varphi(t) = 0$, for all t > 0, which means that $u \equiv v$ in $L^{(p,\infty)}(\mathbb{R}^{N+k})$.

4. Proof of Theorem 1.2

We will split the proof of Theorem 1.2 into small lemmas.

The Lorentz space $L^{(p,q)}(\mathbb{R}^{N+k})$ enjoys the following scaling relation:

(4.1)
$$||f(\lambda \cdot, \lambda^2 \cdot)||_{(p,q)} = \lambda^{-\frac{N+2k}{p}} ||f||_{(p,q)}.$$

To see this, recall that $m(T(E)) = |\det T| m(E)$ for any Lebesgue measurable set E and $T \in GL_n(\mathbb{R})$. Thus, consider the dilation $\delta_{\lambda}(x,y) = (\lambda^{-1}x,\lambda^{-2}y)$, we conclude that

$$m(\delta_{\lambda}(E)) \le t \iff m(E) \le \lambda^{N+2k}t.$$

Using this in the rearrangement function and then in the averaging function (see Section 2), we have

$$(f_{\lambda})^*(t) = f^*(\lambda^{N+2k}t)$$

and

$$(f_{\lambda})^*(t) = f^*(\lambda^{N+2k}t),$$

respectively, where $f_{\lambda} = f(\lambda \cdot, \lambda^2 \cdot)$. Now, this latter identity inside (2.2) implies (4.1). Identity (4.1) says that, if we want the initial condition to comply with (1.15), we must take it in the space $L^{(p,\infty)}(\mathbb{R}^{N+k})$, $p = \frac{N+2k}{2}(\rho-1)$, which is the case of Theorem 1.1. Furthermore, the Marcinkiewcz space $L^{(p,\infty)}(\mathbb{R}^{N+k})$ is invariant by the (1.15):

$$||u_{\lambda}||_{(p,q)} = ||u||_{(p,q)}.$$

Consider the following sequence

(4.2)
$$u_1(x,t) = S_{\mathcal{G}}(t)u_0(x)$$
, and $u_{k+1}(x,t) = u_1(x,t) + B(u_k)$, $k \in \mathbb{N}$.

The proof of Theorem 1.2 relies on the fact that Picard's sequence 4.2 propagates the stated symmetries and positivity of the initial conditions. Indeed, the theorem will be proven if we show that if the initial condition iss the the symmetry (resp. is nonnegative), then each term of Picard's sequence will also be invariant by that symmetry property (resp. is positive). Therefore, it will be the theorem will be proved after the lemmas below.

Lemma 4.1. (a) If φ is a nonnegative function, then $S_{\mathcal{G}}(t)\varphi$ is positive.

(b) Let $\varphi \in L^{(p,\infty)}$ be a homogeneous function of degree $-\frac{2}{\rho-1}$ in the following sense:

(4.3)
$$\varphi(\lambda x, \lambda^2 y) = \lambda^{-\frac{2}{\rho - 1}} \varphi(x, y), \ (x, y) \in \mathbb{R}^{N + k}.$$

Then $S_{\mathcal{G}}(t)\varphi$ is invariant by the scaling map (1.15).

(c) If the function φ is invariant under the action of \mathcal{A} , then $S_{\mathcal{G}}(t)\varphi$ also is. (1.15).

Proof. The proof of (a) is immediate from the positivity of the heat kernel (1.5). To prove (b), we will use expression (1.5):

$$K(\lambda x, \lambda x_{0}, \lambda^{2} y, \lambda^{2} t)$$

$$= \frac{1}{(2\pi)^{\frac{N+2k}{2}}} \int_{\mathbb{R}^{k}} \left(\frac{|\xi|}{\sinh(|\xi|\lambda^{2}t)} \right)^{\frac{N}{2}} e^{i\xi \cdot \lambda^{2} y - \frac{|\xi|}{2} \left((|\lambda x|^{2} + |\lambda x_{0}|^{2}) \coth(|\xi|\lambda^{2}t) - 2(\lambda x) \cdot (\lambda x_{0}) \operatorname{csch}(|\xi|\lambda^{2}t) \right)} d\xi$$

$$= \frac{\lambda^{-2k-N}}{(2\pi)^{\frac{N+2k}{2}}} \int_{\mathbb{R}^{k}} \left(\frac{|\xi|}{\sinh(|\xi|t)} \right)^{\frac{N}{2}} e^{i\xi \cdot y - \frac{|\xi|}{2} \left((|x|^{2} + |x_{0}|^{2}) \coth(|\xi|t) - 2x \cdot x_{0} \operatorname{csch}(|\xi|t) \right)} d\xi$$

$$= \lambda^{-2k-N} K(x, x_{0}, y, t)$$

Then, a change of variables in (1.9), and the homogeneity (4.3) yield

$$(S_{\mathcal{G}}(\cdot)\varphi)_{\lambda}(x,y,t) = \lambda^{\frac{2}{\rho-1}} \int_{\mathbb{R}^{N+k}} K(\lambda x, w, z, \lambda^{2}t) \ \varphi(w, \lambda^{2}y - z) \, \mathrm{d}w \, \mathrm{d}z$$

$$= \lambda^{N+2k+\frac{2}{\rho-1}} \int_{\mathbb{R}^{N+k}} K(\lambda x, \lambda w, \lambda^{2}z, \lambda^{2}t) \ \varphi(\lambda w, \lambda^{2}y - \lambda^{2}z) \, \mathrm{d}w \, \mathrm{d}z$$

$$= (S_{\mathcal{G}}(\cdot)\varphi) (x, y, t).$$

The proof of (c) relies on the fact that $T \in \mathcal{A}$ is written as $T = (T_1, T_2)$ for orthogonal matrices $T_1 \in \mathcal{O}(N)$ and $T_2 \in \mathcal{O}(k)$. Then $|T_1x| = |x|$ and $(T_2^t)^{-1} = T_2$. Hence, the change of variables $T_2^t \xi = \xi'$ gives

$$K(x_0, T(x, y), t) = \frac{1}{(2\pi)^{\frac{N+2k}{2}}} \int_{\mathbb{R}^k} \left(\frac{|\xi|}{\sinh(|\xi|t)} \right)^{\frac{N}{2}} e^{i\xi \cdot T_2 y - \frac{|\xi|}{2} \left((|T_1 x|^2 + |x_0|^2) \coth(|\xi|t) - 2T_1 x \cdot x_0 \operatorname{csch}(|\xi|t) \right)} d\xi$$

$$= K(x, T_1^t x_0, y, t).$$

Then, the change (w, z) = T(w, z) gives $S_{\mathcal{G}}(t)\varphi(T(x, y)) = S_{\mathcal{G}}(t)\varphi(x, y)$.

Lemma 4.2. Let $u \in L^{\infty}(0, \infty; L^{(p,\infty)})$.

- (a) If u is positive, then B(u) is also positive.
- (b) If u is invariant by the scaling map (1.15), then so is B(u).
- (c) If u is invariant under the action of A, then so is B(u).

Proof. The positivity in (a) follows immediately from f(u) > 0 and Lemma 4.1 (a). Notice that, if φ is invariant under the dilation (4.1), then

(4.4)
$$\lambda^{\frac{2\rho}{\rho-1}}\varphi^{\rho}(\lambda x, \lambda^2 y) = (\varphi_{\lambda}(x, y))^{\rho} = \varphi^{\rho}(x, y).$$

Hence,

$$[S_{\mathcal{G}}(\lambda^{2}t)\varphi^{\rho}](\lambda x, \lambda^{2}y) = \int_{\mathbb{R}^{N+k}} K(\lambda x, w, \lambda^{2}y - z, \lambda^{2}t)\varphi^{\rho}(w, z) \, \mathrm{d}w \, \mathrm{d}z$$

$$= \lambda^{N+2k} \int_{\mathbb{R}^{N+k}} K(\lambda x, \lambda w, \lambda^{2}(y - z), \lambda^{2}t)\varphi^{\rho}(\lambda w, \lambda^{2}z) \, \mathrm{d}w \, \mathrm{d}z$$

$$= \int_{\mathbb{R}^{N+k}} K(x, w, y - z, t)\varphi^{\rho}(\lambda w, \lambda^{2}z) \, \mathrm{d}w \, \mathrm{d}z$$

$$= \lambda^{-\frac{2\rho}{\rho-1}} S_{\mathcal{G}}(t)\varphi^{\rho}(x, y)$$

Then,

$$(B(u))_{\lambda} = \left[\int_{0}^{\cdot} S_{\mathcal{G}}(\cdot - s) u^{\rho}(s) ds \right]_{\lambda}$$

$$= \lambda^{\frac{2}{\rho - 1}} \int_{0}^{\lambda^{2} t} S_{\mathcal{G}}(\lambda^{2} t - s) u^{\rho}(\lambda x, \lambda^{2} y, s) ds$$

$$= \lambda^{2 + \frac{2}{\rho - 1}} \int_{0}^{t} S_{\mathcal{G}}(\lambda^{2} (t - s)) u^{\rho}(\lambda x, \lambda^{2} y, \lambda^{2} s) ds$$

$$= \lambda^{2 + \frac{2}{\rho - 1} - \frac{2\rho}{\rho - 1}} \int_{0}^{t} S_{\mathcal{G}}(t - s) u^{\rho}(s) ds$$

$$= \int_{0}^{t} S_{\mathcal{G}}(t - s) u^{\rho}(s) ds$$

$$= B(u).$$

It proves (b). The invariance of f(u) by \mathcal{A} and Lemma 4.1 (c) finish the proof.

5. Final remarks

We finish the paper with some remarks that are either useful to give a more concrete taste to our results or also to give insights to further related research.

- (a) Let T_1 and T_2 be the rotation matrices in \mathbb{R}^N and \mathbb{R}^k , respectively. Then, $T=(T_1,T_2)\in\mathcal{A}$. Then, if $u_0(x,y)=u_0(|x|,|y|)$, Theorem 1.2 guarantees that the solutions are *cylindrical* in the sense that $u(x,y,t)=u(|x|,|y|,t)=u_0(|x|,|y|)$, for all $(x,y,t)\in\mathbb{R}^{N+k}\times(0,\infty)$.
- (b) If $u_0(x,y) = \varepsilon |x|^{-\frac{2}{\rho-1}} |y|^{-\frac{1}{\rho-1}}$, then, for $\varepsilon > 0$ sufficiently small, Theorem 1.2 gives the existence of a positive, self-similar, and symmetric solution of (1.3)–(1.4). Such a solution belongs to $L^{\infty}((0,\infty);L^{(p,\infty)})$.
- (c) Given a initial datum $u_0 \in L^{(p,\infty)}$ satisfying (4.3), the self-similar solution u given by Th. 1.2 is stable in the sense of Eq. (1.12) in Th. 1.1. Nevertheless, the proof of (1.12) tells us that it can be replaced with

(5.1)
$$||u(t) - v(t)||_{(p,\infty)} \le M ||S_{\mathcal{G}}(t) (u_0 - v_0)||_{(p,\infty)}.$$

Therefore, if in addition

$$||S_{\mathcal{G}}(t)(u_0-v_0)||_{(p,\infty)}\to 0,$$

as $t \to \infty$, then

(5.2)
$$\lim_{t \to \infty} ||u(t) - v(t)||_{(p,\infty)} = 0,$$

that is, each self-similar solution is a basin of attraction. Similar results occur with the heat equation with the fractional Laplacian [8] and the Navier-Stokes equation [6].

(d) If a positive classical solution is self-similar and we make $\lambda = t^{-\frac{1}{2}}$ in (1.15), then $u(x,y,t) = t^{-\frac{1}{\rho-1}}v(\xi,\eta)$, where $\xi = t^{-\frac{1}{2}}x$ and $\eta = t^{-1}y$. Straightforward computations show that v must be a solution of

(5.3)
$$\Delta_{\mathcal{G}}v - \nabla_{(\xi,\eta)}v \cdot \left(\frac{\xi}{2},\eta\right) - \frac{1}{\rho-1}v + v^{\rho}.$$

 $\Delta_{\mathcal{G}}$ in (5.3) is our operator (1) in variables (ξ, η) and $\nabla_{(\xi, \eta)}$ is the Euclidean gradient. As

$$\Delta_z v - \nabla_z v \cdot z - \frac{1}{\rho - 1} v + v^{\rho}$$

raised self-interest for being related to positive forward self-similar solutions of $u_t = \Delta u + u^{\rho}$ (see [14] and references thereof), also (5.3) is interesting itself.

(e) The last comment concerns with the nonexistence of global solutions for (1.3)—(1.4). For this, define the following energy functional associated with it:

(5.4)
$$E(u) = \frac{1}{2} \int_{\mathbb{R}^{N+k}} |\nabla_{\mathcal{G}} u|^2 dx - \frac{1}{p+1} \int_{\mathbb{R}^{N+k}} |u|^{p+1} dx.$$

Now, define the Sobolev space

$$H^1_{\mathcal{G}} = \{ u \in L^2(\mathbb{R}^{N+k}) : \nabla_{\mathcal{G}} u \in L^2(\mathbb{R}^{N+k}) \}$$

endowed with the norm

$$||u|| = ||u||_{L^2} + ||\nabla_{\mathcal{G}} u||_{L^2}.$$

Here $\nabla_{\mathcal{G}} = (\nabla_x, |x| \nabla_y)$ denotes the intrinsic gradient of our operator. Let $u_0 \in L^{\infty} \cap H^1_{\mathcal{G}}$ be such that $E(u_0) < 0$. If the solution u belongs to $C((0,T); H^1_{\mathcal{G}})$, then we can apply the following Green identity. Let $B_1(R)$ and $B_2(R)$ be balls of radius R > 0 in \mathbb{R}^N and R^k , respectively. For $B(R) = B_1(R) \times B_2(R) \subset \mathbb{R}^{N+k}$, we have

$$\int_{B(R)} u \Delta_{\mathcal{G}} v = \int_{B_2(R)} \int_{S_1(R)} u \frac{\partial v}{\nu_x} + \int_{B_1(R)} \int_{S_2(R)} v \frac{\partial v}{\nu_x} - \int_{B(R)} \nabla_{\mathcal{G}} u \cdot \nabla_{\mathcal{G}} v.$$

In particular, for $v=u\in C_0^\infty$, for R>0, we get

$$\int_{B(R)} u \Delta_{\mathcal{G}} v = -\int_{B(R)} |\nabla_{\mathcal{G}} u|^2.$$

By density, those identities should work with $v, u \in H^1_{\mathcal{G}}$. Hence, if the solution $u \in C((0,T); H^1_{\mathcal{G}}) \cap L^{\rho+1}$, then we can prove that $E(u(\cdot)) \in C([0,T]) \cap C^1((0,T))$ and

$$\frac{d}{dt}E(u(t)) = -\int_{\mathbb{R}^{N+k}} u_t^2.$$

This can be proved by following [27, Lemma 17.5] with the Green identities above. From this on, the Levine concavity method [23] can be applied to show that $T := T_{\text{max}}(u_0) < \infty$.

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